

Developments in numerical methods (I)

WGNE-meeting, DWD, Offenbach, Germany
24-27 Sept. 2019

Michael Baldauf (DWD)

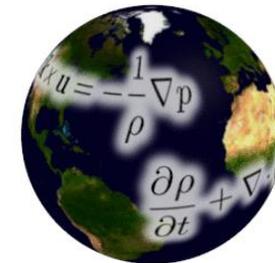
1. report from the ‚PDEs on the sphere 2019‘
2. Planned development at DWD



PDEs on the sphere

29 April – 03 May 2019, Montreal, Kanada

- first workshop 1990, takes place every ~18 months
 - All aspects of dynamical core formulation
 - Coupling between equations and with sub-grid scale parametrisations
 - Parallel scaling
 - Test cases
 - Equation sets
-
- Mostly presentations about meteorology (+ 2* ocean modelling, 2* solar physics)
 - presentations a 20 min. (+ several posters)
 - about 70 participants (16 Can, 12 UK, 11 USA, 11 D, 6 F, ...)
 - <http://collaboration.cmc.ec.gc.ca/science/pdes-2019/index-en.html>



PDEs

THE WORKSHOP ON PARTIAL
DIFFERENTIAL EQUATIONS ON THE SPHERE

Most prominent topics this year

•mixed/compatible finite elements

(talks by Melvin, Lee, Wimmer, Shipton, Gibson, Bendall)

- actual development at UK MO: GungHo/LFRic-project
cubed sphere, better separation between computational and science code;
users desired improved conservation properties

•spectral elements (=continuous Galerkin)

(talks by Reineke, Taylor, Lauritzen)

- US Navy: Operationalisation of NEPTUNE planned for ~2025
(NH, spectral element, cubed sph.)
- E3SM (Energy Exascale Earth System Model) used by 8 DOE labs + univ.
(HOMME-NH dyn core: horiz: SE, HEVI-IMEX)
→ mimetic horizontal differencing → mass & energy conservation

•time integration: **exponential integrators**

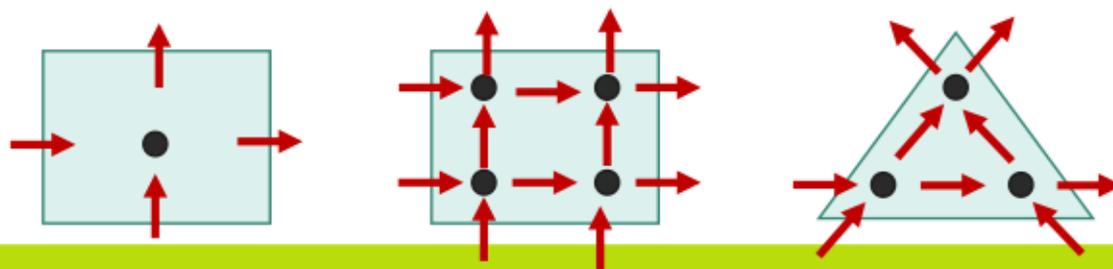
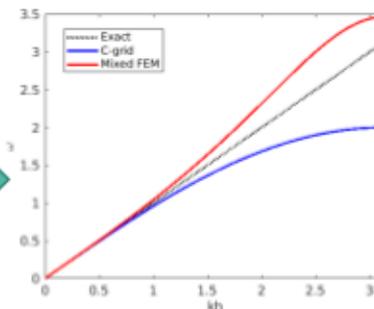
(talks by Côté, Peixoto, Gaudreault)

is a topic since years, but seems not yet used/planned for operational models

Mixed Finite Elements

Mixed Finite Element method gives

- Compatibility: $\nabla \times \nabla \varphi = 0, \nabla \cdot \nabla \times \mathbf{v} = 0$
- Accurate balance and adjustment properties
- No orthogonality constraints on the mesh
- Flexibility of choice mesh (quads, triangles) and accuracy (polynomial order)



Thomas Melvin (UK MO)

Motivation 2 - Exponential integration with nonlinear advection

- Exponential integrators solve accurately **linear** problems
 - Usually allows very large timesteps

$$\frac{dU}{dt} = LU, \quad U(0) = U_0,$$

$$U(t_{n+1}) = e^{\Delta t L} U(t_n),$$

- For nonlinear equations
 - Usually nonlinearity limits timestep sizes

$$\frac{dU}{dt} = LU + N(U), \quad U(0) = U_0,$$

$$U(t_{n+1}) = e^{\Delta t L} U(t_n) + e^{\Delta t L} \int_{t_n}^{t_{n+1}} e^{-(s-t_n)L} N(U(s)) ds,$$

Goal2

How to treat N if dominated by nonlinear advection, in way to allow large time step sizes?

Some other statements:

- P. Lauritzen: total energy budgets
 - „Conserving total energy to within $\sim 0.01 \text{ W/m}^2$ is considered ‚good enough‘ for coupled climate modeling“
compare: the earths energy imbalance is $\sim 1 \text{ W/m}^2$
 - In CAM-SE: total energy errors in the dyn. core, physics-dynamics coupling and pressure work errors are $\sim -0.6 \dots +0.3 \text{ W/m}^2$.
Local errors can be an order of magnitude larger (at least).
- Xi Chen:
grid-staggering becomes increasingly less important for higher-order methods

A possible alternative dynamical core for ICON based on Discontinuous Galerkin Discretisation

Michael Baldauf (Deutscher Wetterdienst)

MetStröm



Discontinuous Galerkin (DG) methods in a nutshell

$$\frac{\partial q^{(k)}}{\partial t} + \nabla \cdot \mathbf{f}^{(k)}(q) = S^{(k)}(q), \quad k = 1, \dots, K$$

weak formulation

$$\int_{\Omega_j} dx v(\mathbf{x}) \cdot \dots$$

Finite-element ingredient

$$q^{(k)}(x, t) = \sum_{l=0}^p q_{j,l}^{(k)}(t) p_l(x - x_j)$$

e.g. Legendre-Polynomials

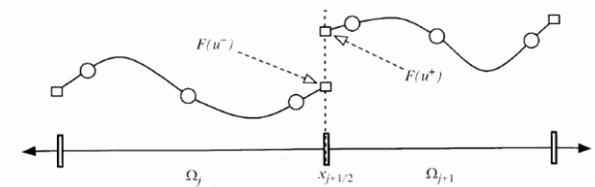
Finite-volume ingredient

$$\mathbf{f}(q) \rightarrow \mathbf{f}^{num}(q^+, q^-) = \frac{1}{2} (\mathbf{f}(q^+) + \mathbf{f}(q^-) - \alpha(q^+ - q^-))$$

Lax-Friedrichs flux

Gaussian quadrature for the integrals of the weak formulation

→ ODE-system for $q^{(k)}_{jl}$



From Nair et al. (2011) in
'Numerical techniques for global atm.
models'

e.g.

Cockburn, Shu (1989) *Math. Comput.*

Cockburn et al. (1989) *JCP*

Hesthaven, Warburton (2008):
Nodal DG Methods

- **local conservation**
- any **order of convergence** possible
- flexible application on **unstructured grids** (also dynamic adaptation is possible, h-/p-adaptivity)
- very good **scalability**
- **explicit** schemes are easy to build and are quite well understood
- higher accuracy helps to **avoid several awkward approaches** of standard 2nd order schemes: staggered grids (on triangles/hexagons, vertically heavily stretched), numerical hydrostatic balancing, grid imprints by pentagon points or along cubed sphere lines, ...

- high computational costs due to
 - (apparently) **small Courant numbers**
 - higher number of DOFs
- **well-balancing** (hydrostatic, perhaps also geostrophic?) in Euler equations is an issue (can be solved!)
- basically ,only‘ an **A-grid-method**, however, the ,spurious pressure mode‘ is very selectively damped!

Target system: ICON model

(Zängl et al. (2015) QJRMS)

- operational at DWD since Jan. 2015
(global (13km) and nest over Europe (6.5km))
- convection-permitting (2.2km): Q4/2020
- horiz.: icosahedral triangle C-grid, vertic.: Lorenz-grid
- non-hydrostatic, compressible
- mixed finite-volume / finite-difference (mass, tracer mass conservation)
- predictor-corrector time-integration → overall 2nd order discretization

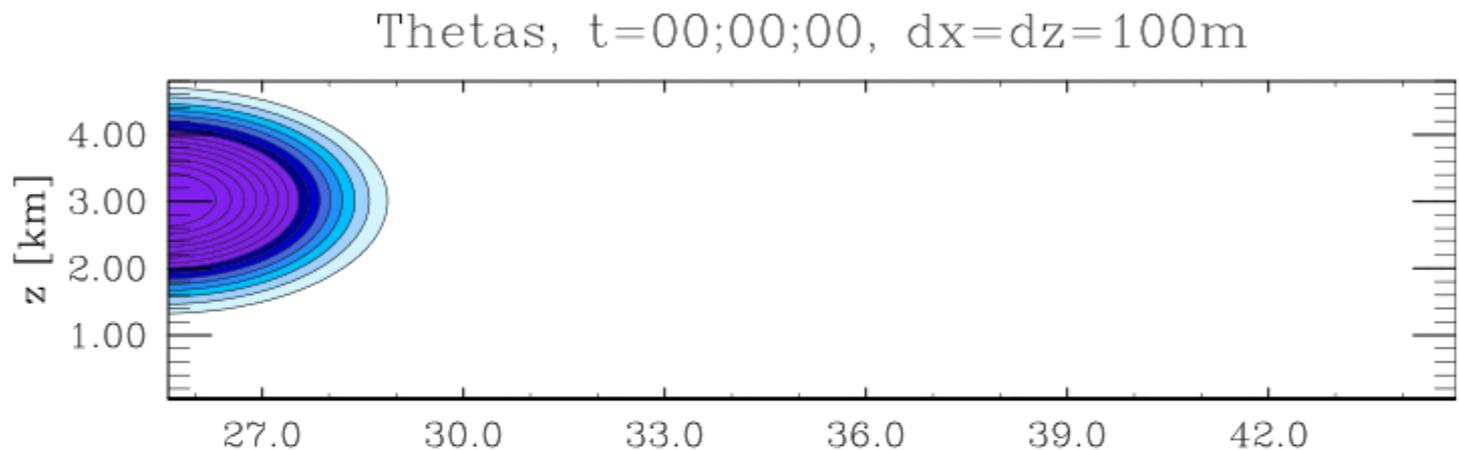


but currently far away from this, only a **toy model for 2D problems exists** with:

- explicit time integration DG-RK (with Runge-Kutta schemes) or horizontally explicit-vertically implicit (DG-HEVI) (with IMEX-Runge-Kutta)
- ‚local DG‘ (LDG) option for PDEs with higher spatial derivatives
- use of a triangle grid (also on the sphere) is optional

Test case: falling cold bubble

Testsetup by *Straka et al (1993)*



Test properties:

- test of dry Euler equations (without Coriolis force)
- unstationary
- strongly nonlinear
- comparison with reference solution from paper

COSMO

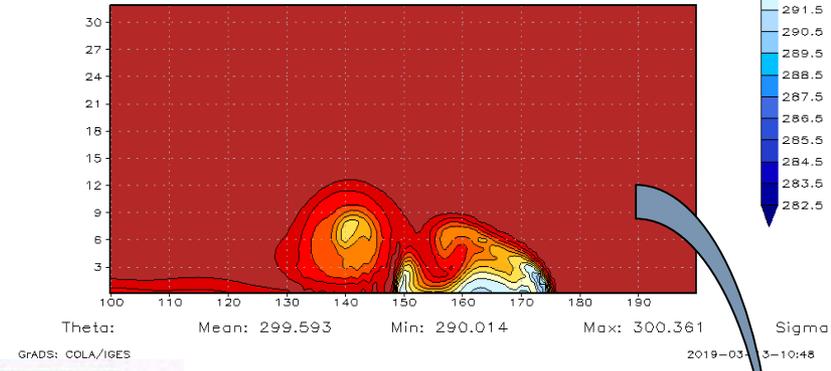
RK2-SLC
dt=0.0826446280992
dx=200.0
dy=200.0

DG explicit

2nd order

dx=dz=200m

Theta p=1 t=900.0

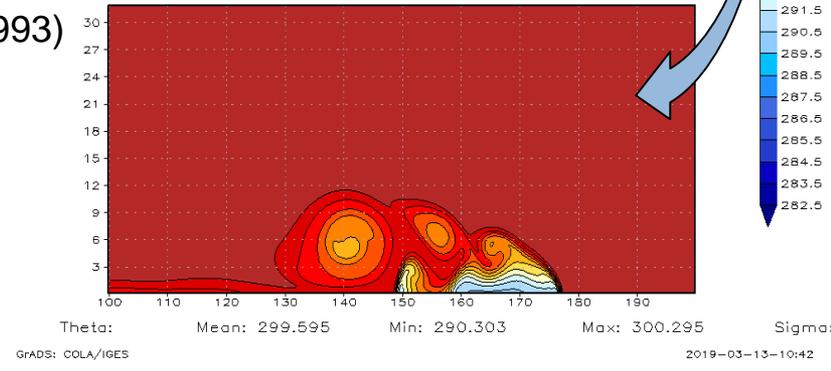


Faktor 4.3
in comput. time

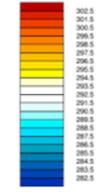
3rd order

dx=dz=200m

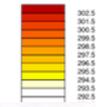
Theta p=2 t=900.0



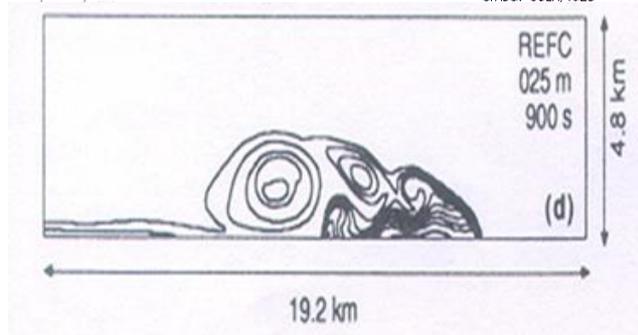
Theta (in K)



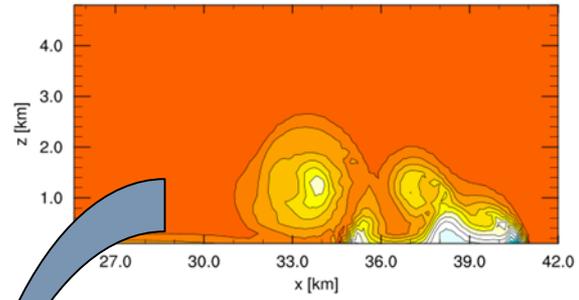
Theta (in K)



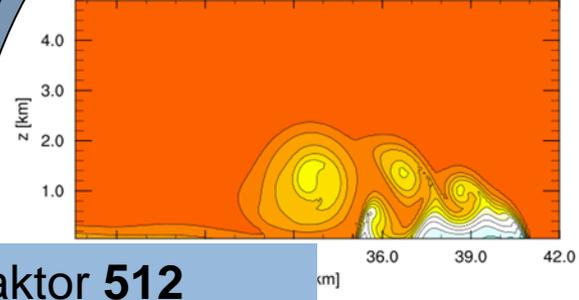
Reference solution
from Straka et al. (1993)



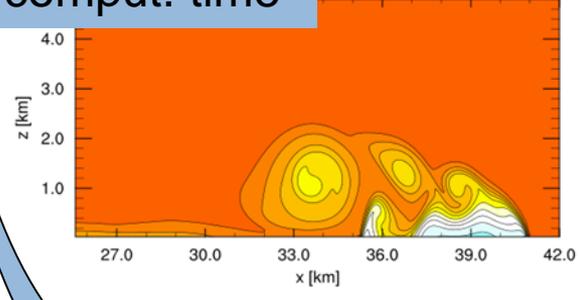
t=00:15:00, dx=dz=200m,



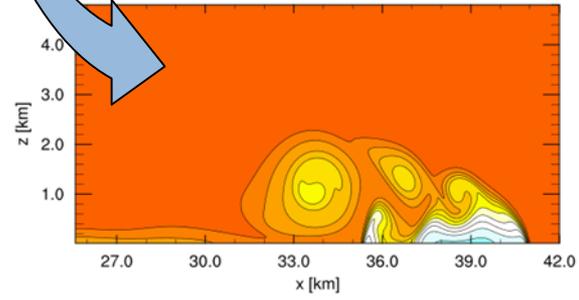
t=00:15:00, dx=dz=100m,



t=00:15:00, dx=dz=50m,



t=00:15:00, dx=dz=25m,

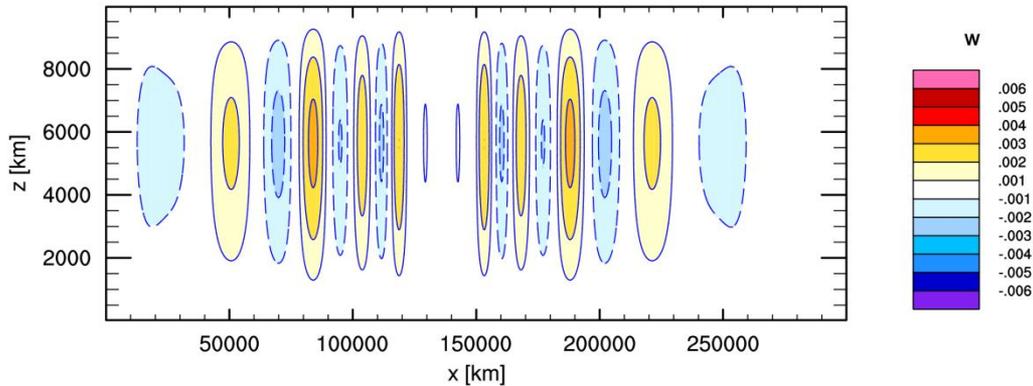


Faktor 512
in comput. time

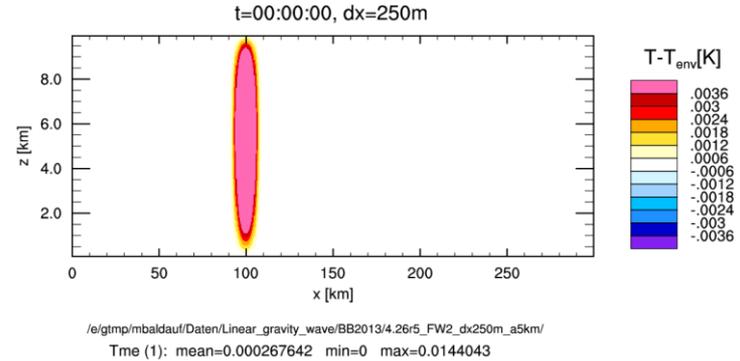
setup similar to *Skamarock, Klemp (1994) MWR*

$\Delta x=500\text{m}, \Delta z=250\text{m}$

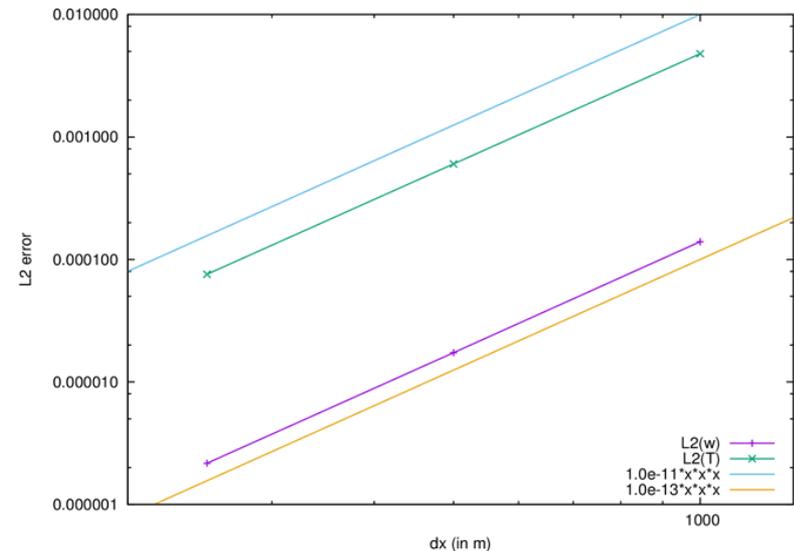
w, t=00:30:00



colors : simulation with $p=2/\text{RK3-SSP}$ (i.e. 3rd order explicit DG)
blue lines: analytic solution for compressible, non-hydrostatic Euler eqns. (*Baldauf, Brdar (2013) QJRMS*)



Exact 3rd order convergence for w and T':



Horizontally explicit - vertically implicit (HEVI)-scheme with DG

Motivation: get rid of the strong time step restriction by vertical sound wave expansion in flat grid cells (in particular near the ground)

$$\frac{\partial q^{(s)}}{\partial t} + \underbrace{\nabla \cdot \mathbf{f}_{slow}^{(s)}}_{\text{explicit}} + \underbrace{\nabla \cdot \mathbf{f}_{fast}^{(s)}}_{\text{implicit}} = \underbrace{S_{slow}^{(s)}}_{\text{explicit}} + \underbrace{S_{fast}^{(s)}}_{\text{implicit}} \quad \mathbf{f}_{fast}^{(s)} = f_{z,fast}^{(s)} \mathbf{e}_z$$

$$f_{z,fast}^{(s)} = \sum_{s'} H^{ss'} q^{(s')}$$

- Use of IMEX-RK (SDIRK) schemes: SSP3(3,3,2), SSP3(4,3,3) (*Pareschi, Russo (2005) JSC*)
- The implicit part leads to several band diagonal matrices
→ here a direct solver is used (expensive!)

References:

Giraldo et al. (2010) SIAM JSC: propose a HEVI semi-implicit scheme

Bao, Klöfkorn, Nair (2015) MWR: use of an iterative solver for HEVI-DG

Blaise et al. (2016) IJNMF: use of IMEX-RK schemes in HEVI-DG

Abdi et al. (2017) arXiv: use of multi-step or multi-stage IMEX for HEVI-DG

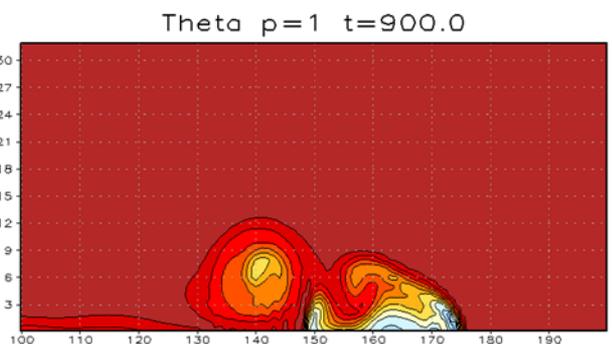
Test case: falling cold bubble (Straka et al. (1993))

Comparison explicit vs. HEVI scheme

2nd order

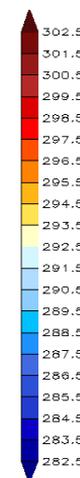
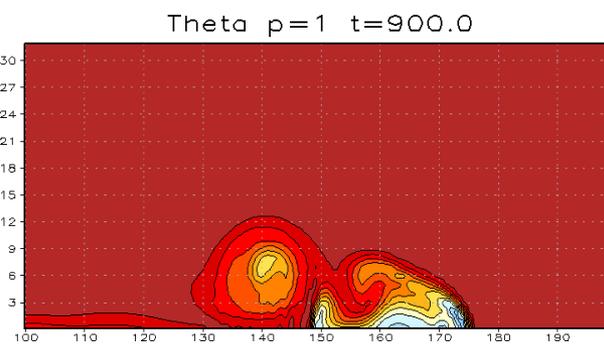
RK2-SLC
dt=0.0826446280992
dx=200.0
dy=200.0

DG explicit

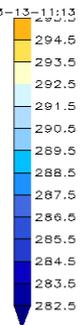
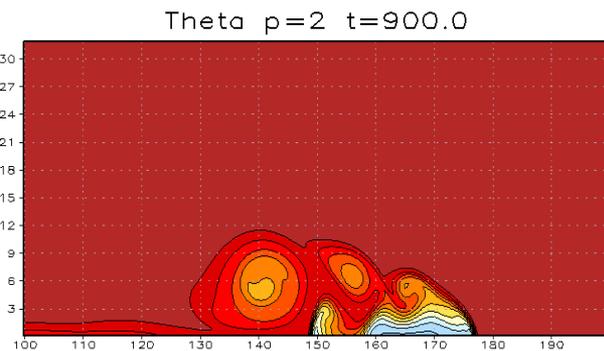
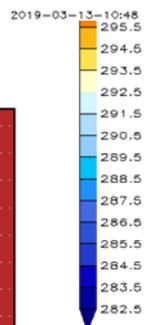
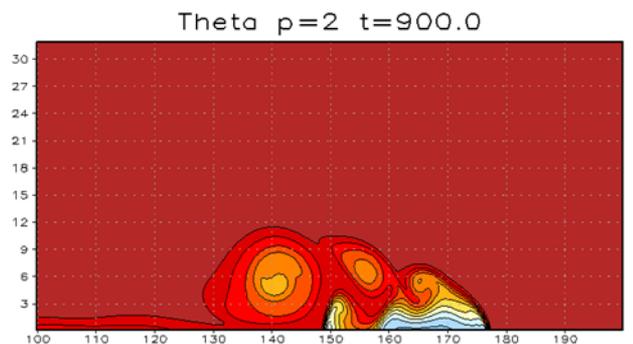


SSP3-3-3-2
dt=0.16393442623
dx=200.0
dy=200.0

DG HEVI



3rd order



How to bring DG on the sphere ...

Idea to avoid pole problem and to keep high order discretization:
use **local (rotated) coordinates** for every (triangle) grid cell,
i.e. rotate every grid cell towards $\lambda \approx 0$, $\varphi \approx 0$.

→ geometry is treated exactly

→ transform fluxes between neighbouring cells

shallow water equations

covariant formulation (here: without bathymetry)

$$\frac{\partial \sqrt{GH}}{\partial t} + \frac{\partial}{\partial x^i} \sqrt{G} m^i = 0$$

$$\frac{\partial \sqrt{G} m^i}{\partial t} + \frac{\partial}{\partial x^j} \sqrt{G} T^{ij} = \sqrt{G} (F_{Cor}^i - \Gamma_{jk}^i T^{jk})$$

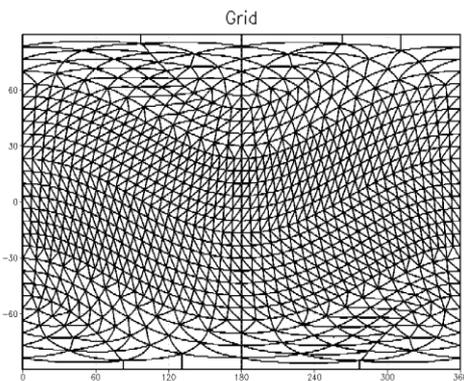
$$T^{ij} = \frac{m^i m^j}{H} + \frac{1}{2} g^{ij} g_{grav} H^2$$



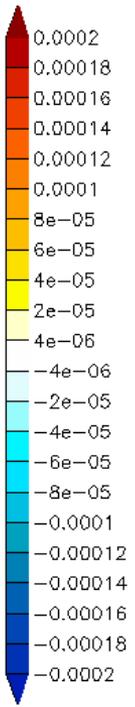
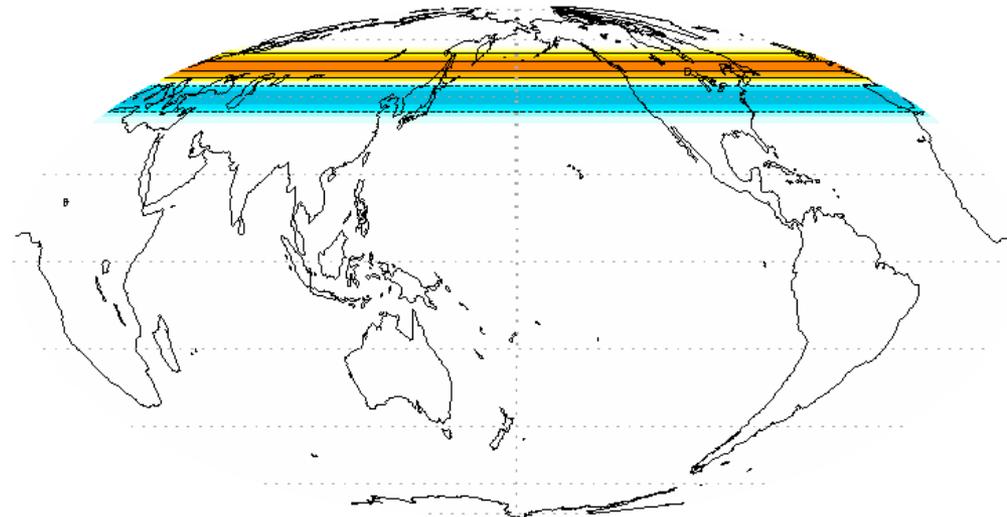
Barotropic instability test *Galewsky et al. (2004)*

4th order DG scheme
without additional diffusion
 $dx \sim 67$ km, $dt = 15$ sec.

simple triangle grid
on the sphere
 $dx \sim 500$ km:



rel. Vortic., ord=4 t=0d00h00m0.0s



relVort:

Mean: 7.53016e-07 Min: -9.8335e-05 Max: 0.000112421 Sigma: 2

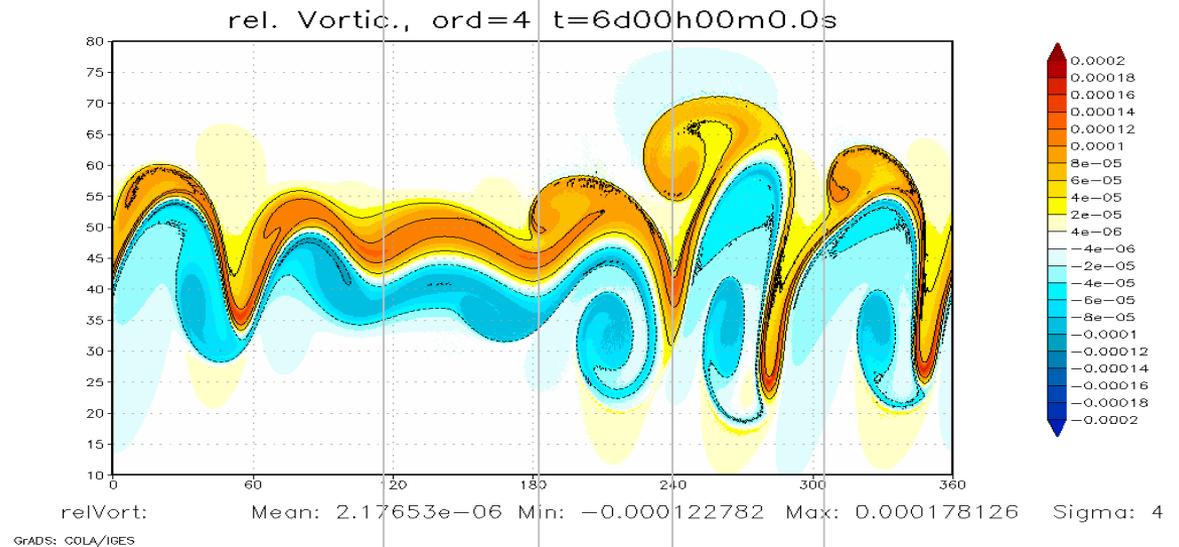


Barotropic instability test

Galewsky et al. (2004)

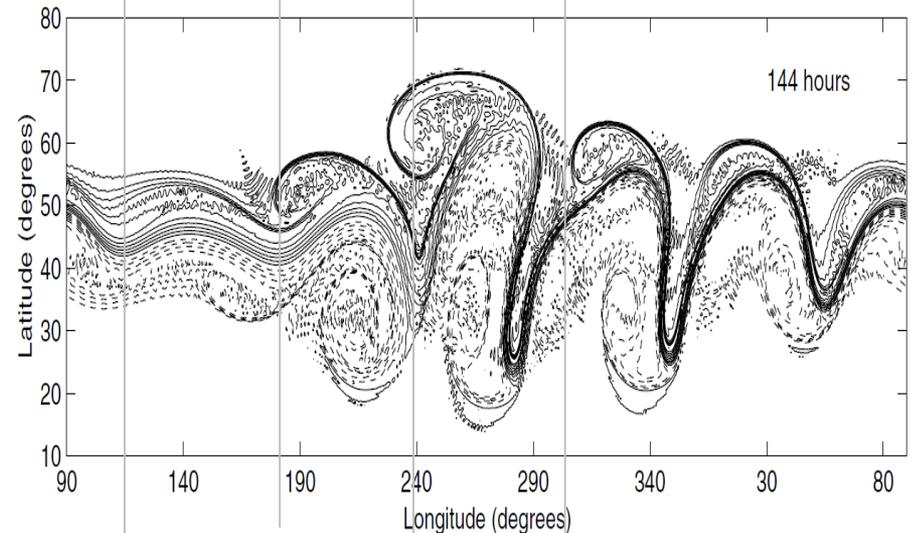
4th order DG scheme
without additional diffusion
dx~67 km, dt=15 sec.

relative vorticity



FMS-SWM (Geophys. Fl. Dyn. Lab.)
without additional diffusion
dx~60 km (T341), dt=30 sec.

Fig. 4 from *Galewsky et al. (2004)*

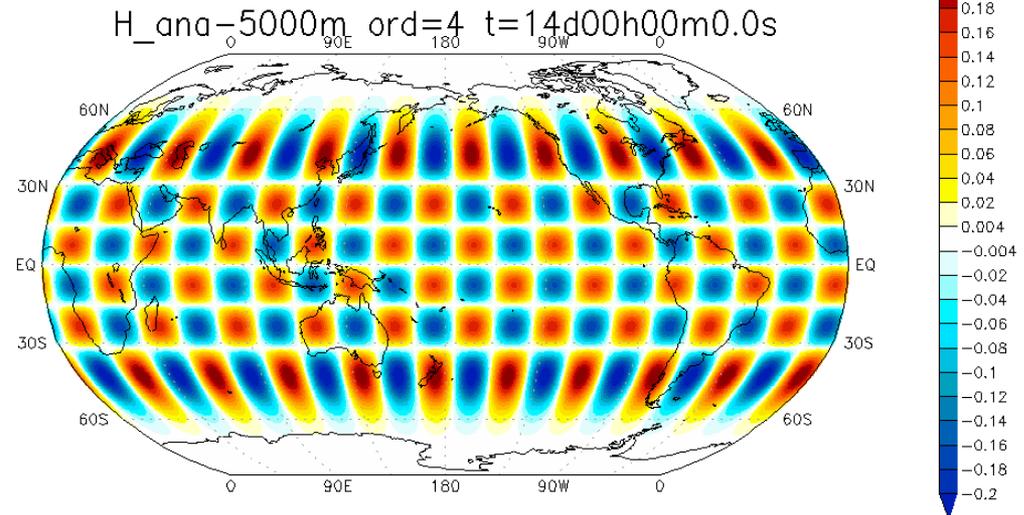
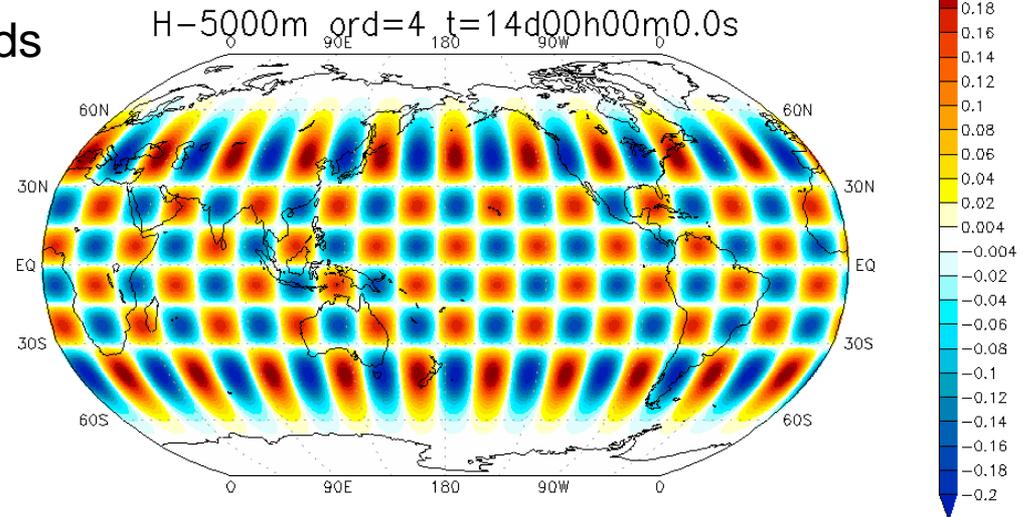


Linear inertial-gravity wave

Solution after 14 days ~ 100 periods

4th order DG scheme
without additional diffusion
dx~134 km, dt=30 sec.

→ wave pattern remains stable



Analytic solution by
Shamir, Paldor (2016),
Paldor (2013)

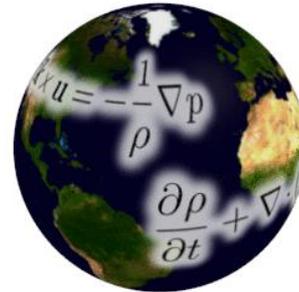
(is known to be slightly too fast
by ~0.3% ~0.3 wavelengths)
This test case might possibly replace the
well-known 'Rossby-Haurwitz wave test'
in the future.

Summary

- **2D toy model** for
 - explicit DG-RK (on arbitrary unstructured grids with triangle or quadrilateral grid cells) and
 - **HEVI DG-IMEX-RK**
works for several idealized tests (also Euler equations with terrain-following coordinates), correct convergence behaviour, ...
- **DG on the sphere** by use of local (rotated gnomonial) coordinates

Outlook

- **further design decisions:** nodal vs. modal, local DG vs. interior penalty vs. ..., ..
- coupling of **tracer advection** (mass-consistency)?
- improve **efficiency** in the HEVI direct solver
- further **milestones** (for the next years!)
 - development of a 3D prototype DG-HEVI solver
 - choose optimal convergence order p and grid spacing
estimated: $p_{\text{horiz}} \sim 3 \dots 6$, $p_{\text{vert}} \sim 3 \dots 4$ ($p_{\text{time}} \sim 3 \dots 4$)



PDEs

THE WORKSHOP ON PARTIAL
DIFFERENTIAL EQUATIONS ON THE SPHERE

Announcement:

The next

„Partial differential equations on the sphere“ – workshop

will take place at
DWD, Offenbach, Germany
5-9 October 2020